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# Association probabilities between the single chain amphiphiles into a binary mixture in planar monolayers

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In this work, we have calculated the probabilities of selfassociation of single chain amphiphile molecules for 36 binary mixtures of molecules having an even number of carbon atoms (from 8 to 24 carbon atoms). The association probabilities depend on the dipole moment of head groups, the length of chains and the mole fraction of the mixture of both types of amphiphile. These calculations are useful for the quantitative and qualitative explanation of the cluster formation on the surface of the mono- or bilayers.

## Introduction

Owing to the fact that in natural membranes there are different types of lipids, useful information can be obtained studying supermolecular aggregates made up of mixtures of two or more species of lipid molecules. Because the lipids in mixtures do not chemically react, the supermolecular aggregates may be considered as athermal solutions. The distribution of the components into the structure of the supermolecular aggregate is not done at random, but it is due to a selective association process among lipid molecules in the mixture. Jain [1] wrote about the existence of the microclusters separated by the microinterfaces in the mixtures of synthetic lipids. In this work we have calculated the selfassociation probabilities between single chain amphiphiles into binary mixtures organized as a planar monolayer.

## Association probabilities

If we want to obtain the expressions of the association probabilities, we assume an athermal solution of the single hydrocarbon chain molecules having different radii,  $r$ , and dipole moments,  $\bar{p}$ , and also the chains length, mixed up in the mole fractions  $x_1$  and  $x_2$  and organized in a planar monolayer or bilayer. It is clear

that the selectivity of association of a molecule to another molecule of the same type, or of different type is determined by the difference between the coupling energies, which are determined by the interaction forces between the associating molecules. Evidently, the molecules belonging to a couple interact through the electric dipoles of the polar head groups and through the van der Waals-London dispersion forces between the hydrocarbon chains. The interaction energies between two molecules of the same type,  $U_{11}$  and  $U_{22}$  or different type,  $U_{12}$ , have unequal values, and determine a selective association whose result is the formation of the clusters on the surface of the aggregate. We introduce the probabilities of formation of some couples  $P_{11}$ ,  $P_{12}$ ,  $P_{22}$  between which there are the following relations:

$$P_{11}/P_{12} = (x_1/x_2) \cdot \exp(-(U_{12} - U_{11})/RT) \quad (1)$$

$$P_{22}/P_{12} = (x_1/x_2) \cdot \exp(-(U_{12} - U_{22})/RT) \quad (2)$$

$$P_{11} + P_{12} + P_{22} = 1 \quad (3)$$

Solving this system, we have obtained the relations out of which the association probabilities are calculated:

$$P_{11} = x_1^2 e_1 / (x_1^2 e_1 + x_2^2 e_2 + x_1 x_2) \quad (4)$$

$$P_{12} = x_1 x_2 / (x_1^2 e_1 + x_2^2 e_2 + x_1 x_2) \quad (5)$$

$$P_{22} = x_2^2 e_2 / (x_1^2 e_1 + x_2^2 e_2 + x_1 x_2) \quad (6)$$

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with following notations:

$$e_1 = \exp(-(U_{12} - U_{11})/RT); \quad e_2 = \exp(-(U_{12} - U_{22})/RT) \quad (7)$$

$$U_{11} = -(p_1^2 N_A / (4\pi\epsilon(2r_1)^3) + W_{11}) \quad (8)$$

$$U_{12} = -(p_1 p_2 N_A / (4\pi\epsilon(r_1 + r_2)^3) + W_{12}) \quad (9)$$

$$U_{22} = -(p_2^2 N_A / (4\pi\epsilon(2r_2)^3) + W_{22}) \quad (10)$$

where,  $\vec{p}_1$ ,  $r_1$  and  $\vec{p}_2$ ,  $r_2$  are the electric dipole moment and the radius of the polar head group of a molecule in the first, respectively, in the second species of mixture.  $W_{11}$ ,  $W_{12}$ ,  $W_{22}$  are the molal energies of interaction through van der Waals-London dispersion forces between the chains of two molecules of the same type ( $W_{11}$ ,  $W_{22}$ ) or of a different type ( $W_{12}$ ).

In the calculation of the dispersion energy we consider that in the saturated paraffin chains the electronic charge distribution is localized in bonds and atoms. Also the electronic correlation is negligible between distant points in one molecule and the covalent bonds of each carbon atom have a such spatial arrangement so that we can suppose that a  $\text{CH}_2$  group is an attractive center. Under these conditions, within the limits of second-order perturbation theory it is acceptable to assume that dispersion forces are locally additive. So, the attraction force between two hydrophobic chains is obtained by adding all the forces between pairs of attractive centers one in each molecule [2].

The distance between the nearest attractive centers is about 7.14 Å, equal to the distance between the axes of two neighbours molecules with short chains. This is large compared with the dimensions of the attraction units (atoms, bonds), such as the interactions between them, verify London's invers-sixth power distance law ( $W = -A/d^6$ ), where  $A$  is equal to  $5.6 \cdot 10^3 \text{ Å}^6 \cdot \text{J} \cdot \text{kmol}^{-1}$  [2].

Because the distance between the most shorter than 1000 Å, the retardation effects are not taken into account [3]. Owing to the fact that in the calculation of the selfassociation probabilities appears the difference between interaction energies, we think that the non-ad-

ditivity of the dispersion forces does not modify our results and conclusions.

Selfassociation probabilities of the amphiphile molecules in a binary mixture depend on seven parameters:

- (i) the lengths of the hydrocarbonic chains, identified with their number of carbon atoms  $N_1$  and  $N_2$ ;
- (ii) the electric dipole moments of polar head groups  $\vec{p}_1$  and  $\vec{p}_2$ ;
- (iii) the area per amphiphile occupied by the head group of each of the two components  $a_1$  and  $a_2$ ;
- (iv) the molal fraction of one of the two components  $x_1$  ( $x_2 = 1 - x_1$ )

We have imagined 36 binary systems in order to study the dependence of association probabilities of the two components on their chains length. These binary systems were arranged in eight series (sets) of mixtures (Table I). The head of series has the hydrocarbon chain with an even number of carbon atoms  $N_1$  over the range 8–22. The second component has the hydrocarbonic chain with an even number of carbon atoms  $N_2$  ranging between  $N_1$  and 24. The first component has the surface area and the electric dipole moment of the polar head group equal to 40 Å<sup>2</sup> [4] and 5 D, respectively, and the second component has the area equal to 60 Å<sup>2</sup>. These three parameters ( $a_1$ ,  $a_2$  and  $p_1$ ) were considered fixed and available for all the samples.

The last two parameters ( $p_2$  and  $x_1$ ) were modified as follows: for each sample the dipole moment of the second component is equal to 5, 10, 15, 20, 25, 30, 35, 40, 45 D and for each value of the dipole moment ( $p_2$ ), the molal fraction ( $x_1$ ) of the first component is equal to 0.1; 0.2; ... 0.9. Therefore an enormous number of data for association probabilities was calculated. This paper will described only as much as necessary in order to see the influence of each of the variable parameters on the association probabilities.

For an easy rendering the following notations will be used for the association probabilities:  $P_{11}(a, a/a, b)$ ,  $P_{12}(a, a/a, b)$ ,  $P_{22}(b, b/a, b)$  in which the number and letters have the meaning presented in Table I. The index 1 refers to the component with the shorter chain (also named in the text 'the head of the series' or 'the first

TABLE I

*The composition of the single chain lipid mixtures for which the association probabilities in planar bilayers has been calculated*

| The serie | The first component: $a$ | The second component: $b$   |
|-----------|--------------------------|---|
| I         | C <sub>8</sub>           | C <sub>10</sub> ; C <sub>12</sub> ; C <sub>14</sub> ; C <sub>16</sub> ; C <sub>18</sub> ; C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(8,8/8,b)$ $P_{22}(b,b/8,b)$ $P_{12}(8,b/8,b)$ |
| II        | C <sub>10</sub>          | C <sub>12</sub> ; C <sub>14</sub> ; C <sub>16</sub> ; C <sub>18</sub> ; C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(10,10/10,b)$ $P_{22}(b,b/10,b)$ $P_{12}(10,b/10,b)$             |
| III       | C <sub>12</sub>          | C <sub>14</sub> ; C <sub>16</sub> ; C <sub>18</sub> ; C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(12,12/12,b)$ $P_{22}(b,b/12,b)$ $P_{12}(12,b/12,b)$                               |
| IV        | C <sub>14</sub>          | C <sub>16</sub> ; C <sub>18</sub> ; C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(14,14/14,b)$ $P_{22}(b,b/14,b)$ $P_{12}(14,b/14,b)$   |
| V         | C <sub>16</sub>          | C <sub>18</sub> ; C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(16,16/16,b)$ $P_{22}(b,b/16,b)$ $P_{12}(16,b/16,b)$   |
| VI        | C <sub>18</sub>          | C <sub>20</sub> ; C <sub>22</sub> ; C <sub>24</sub> $P_{11}(18,18/18,b)$ $P_{22}(b,b/18,b)$ $P_{12}(18,b/18,b)$   |
| VII       | C <sub>20</sub>          | C <sub>22</sub> ; C <sub>24</sub> $P_{11}(20,20/20,b)$ $P_{22}(b,b/20,b)$ $P_{12}(20,b/20,b)$   |
| VIII      | C <sub>22</sub>          | C <sub>24</sub> $P_{11}(22,22/22,b)$ $P_{22}(b,b/22,b)$ $P_{12}(22,b/22,b)$   |

component') and the index 2 refers to the component with the longer chain (also named in the text 'the second component'). Thus,  $P_{11}$  is the direct association

probability of two molecules with the short chain,  $P_{22}$  is the direct association probability of two molecules with the long chain and  $P_{12}$  is the cross association probab-

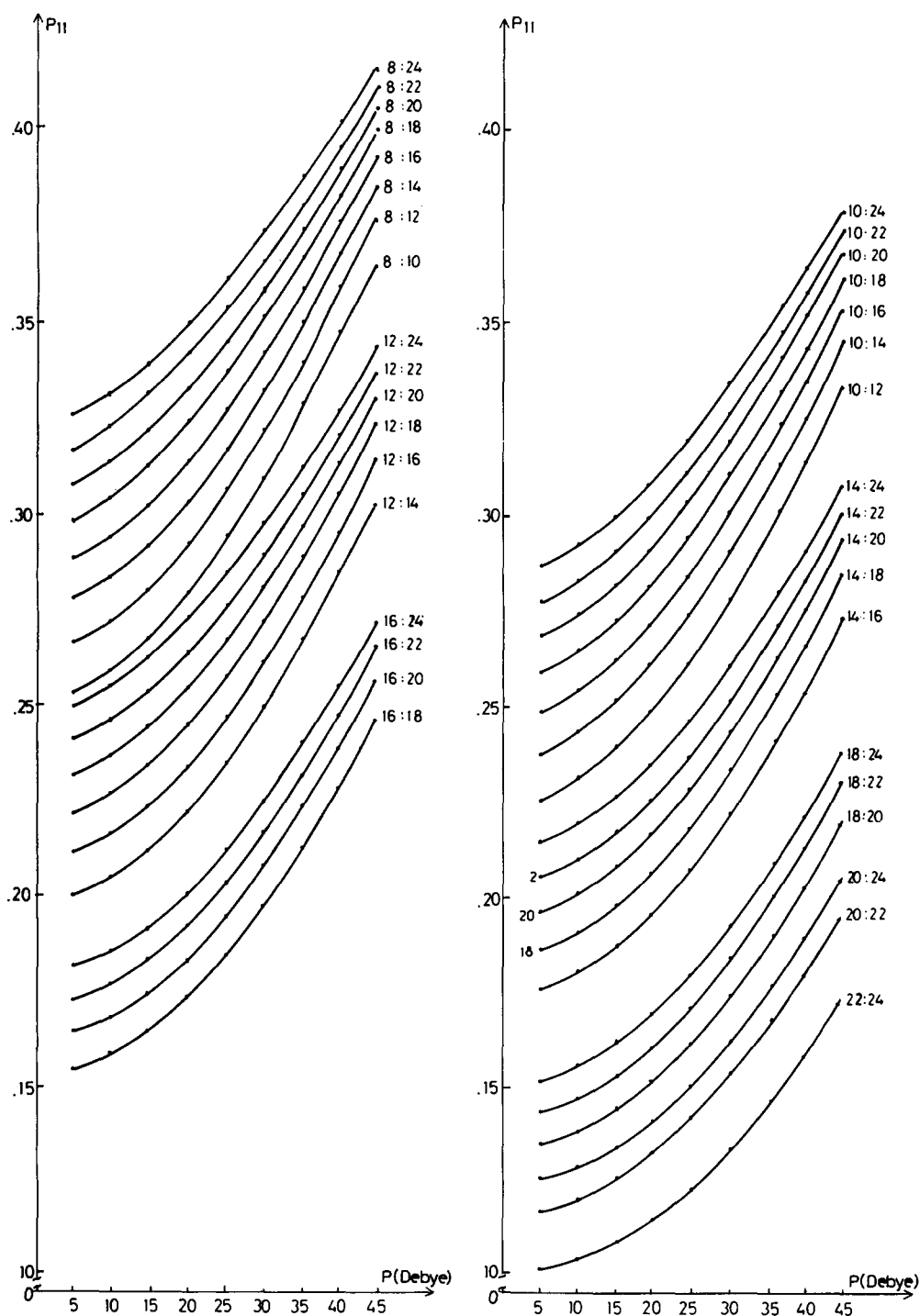
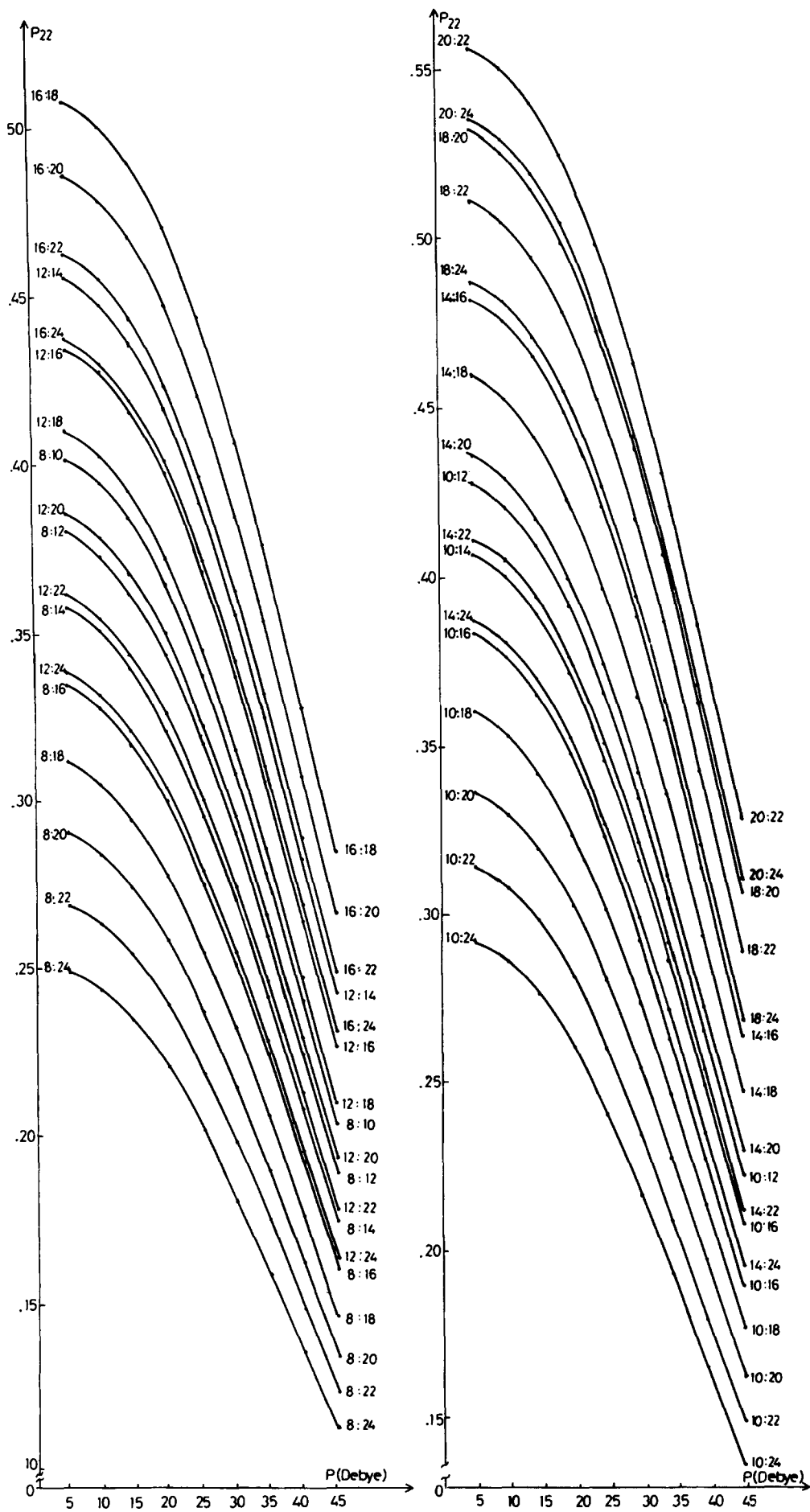


Fig. 1. The variation of selfassociation probability of two short hydrophobic chain molecules,  $P_{11}$ , in a binary mixture depending on the electric dipole moment of the second component. The two components do not chemically react and  $x_1 = x_2 = 0.5$ .

Fig. 2. The variation of selfassociation probability of two long hydrophobic chain molecules,  $P_{22}$ , in a binary mixture depending on the electric dipole moment of the second component.  $x_1 = x_2 = 0.5$ .



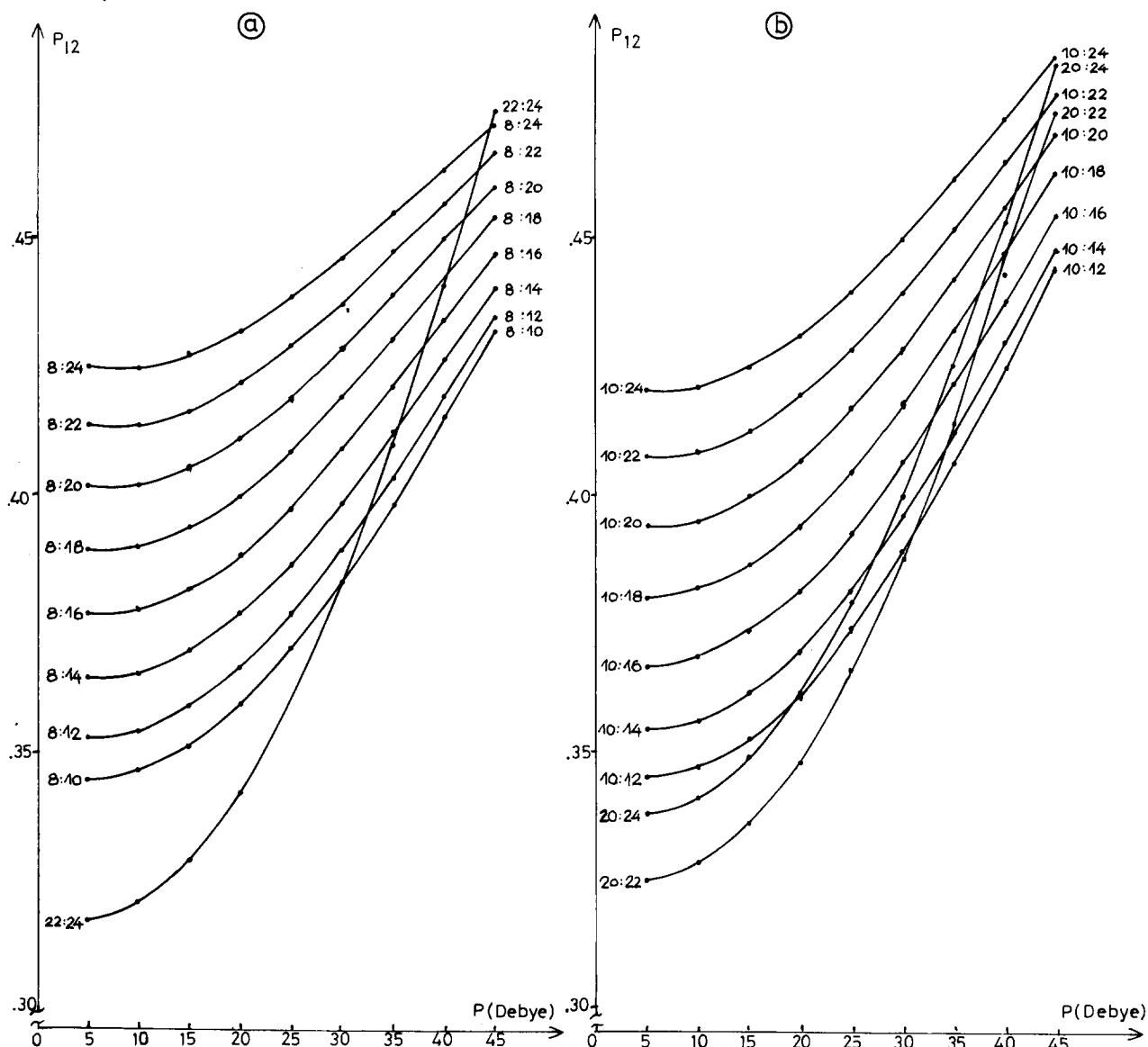


Fig. 3. The variation of cross association probability  $P_{12}$  depending on the electric dipole moment of the second component in a binary mixture in which the components do not chemically react. The series in which the first component with constant dipole moment (5 D) has 8 and 22 carbon atoms (a); 10 and 20 carbon atoms (b); 12 and 18 carbon atoms (c); 14 and 16 carbon atoms (d). In the Figs. 1–3 at the ends of each curve the mixture for which it was drawn was marked. The two numbers represent the chains length (the number of carbon atoms) of the two components in the mixture. The first number identifies the head of the series, the second regards the species with a variable dipole moment. Also,  $x_1 = x_2 = 0.5$ .

ity of two molecules with unequal chains, into a binary mixture. The letters to the left of the line (/) represent the chain length (the number of carbon atoms) of the association molecules and those to the right represent the chain length,  $N_1$  and  $N_2$ , of molecules the mixture was made up of.

## Results and Discussions

Looking at Figs. 1–3, we can easily see that regardless of the difference between the lengths of the chains of the two components,  $P_{11}$  and  $P_{12}$  increases and  $P_{22}$  decreases as the dipole moment of the long chain mole-

cules increases. In Fig. 1, some groups of curves may be noticed, each group corresponding to a series of binary mixtures, in which the first component is the same. As one can easily see, the curves from each group move towards the top of the picture as the second number increases.

On the other hand the curves corresponding to those binary mixtures (comprised in a vertical row of Fig. 9) in which the second component has a certain chain length, move towards the top of the picture as the first number decreases.

In other words, for each value of the electric dipole moment  $p_2$ , the association probability of the shorter

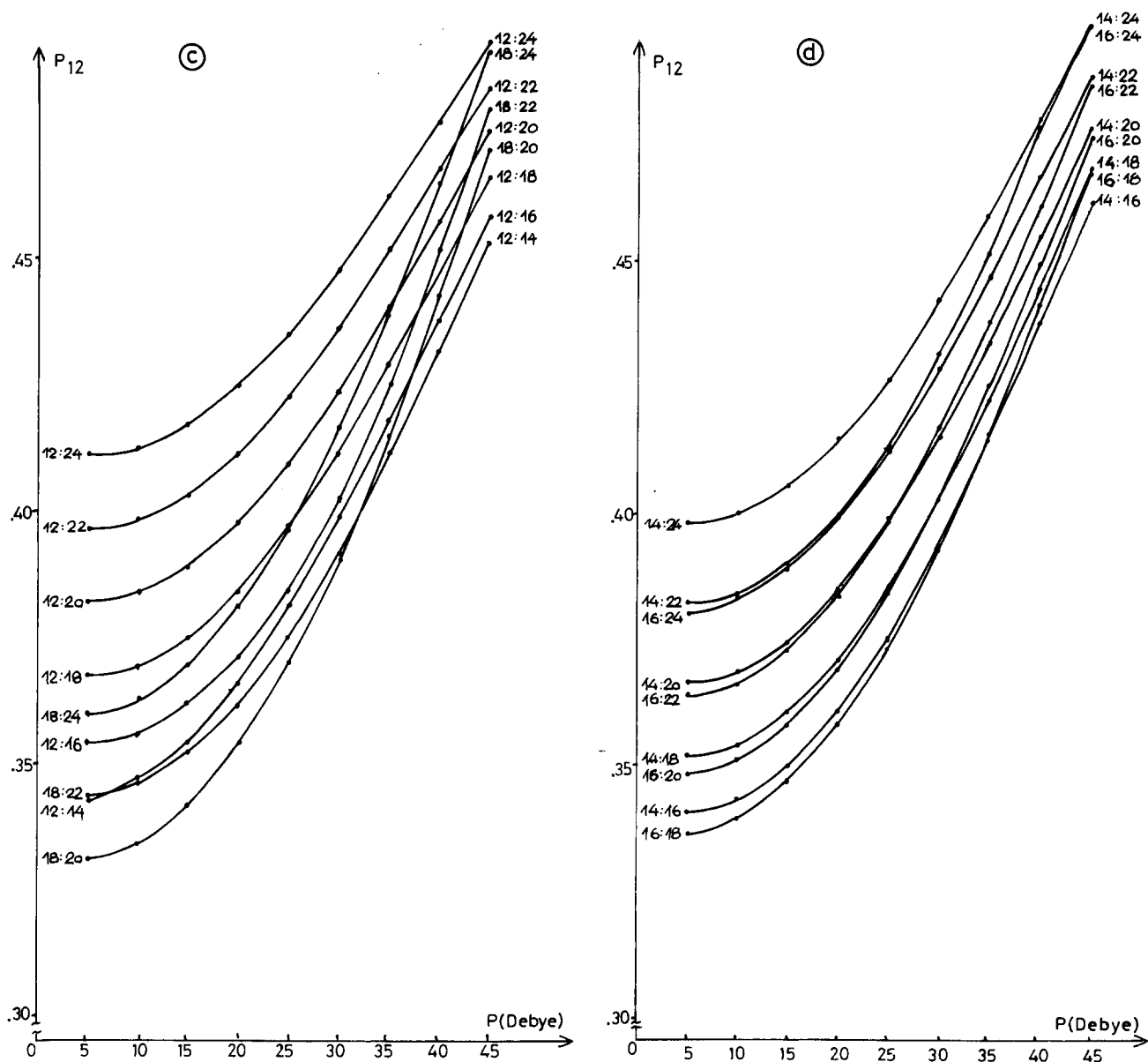


Fig. 3 (continued).

chain molecules  $P_{11}$  increases when the chain length,  $N_2$ , of the second component increases and  $N_1$  is assumed constant, but decreases when the chain length of the first component,  $N_1$ , increases and  $N_2$  is maintained constant (Fig. 1). So, we can write the following table of inequalities for the selfassociation probability  $P_{11}$ :

$$\begin{array}{c}
 P_{11}(8,8/8,10) < P_{11}(8,8/8,12) < P_{11}(8,8/8,14) < \dots < P_{11}(8,8/8,24) \\
 \quad \quad \quad \downarrow \quad \quad \quad \downarrow \quad \quad \quad \downarrow \\
 P_{11}(10,10/10,12) < P_{11}(10,10/10,14) < \dots < P_{11}(10,10/10,24) \\
 \quad \quad \quad \downarrow \\
 P_{11}(12,12/12,14) < \dots < P_{11}(12,12/12,24) \\
 \quad \quad \quad \downarrow \\
 \dots \dots \dots \downarrow \\
 \quad \quad \quad P_{11}(22,22/22,24)
 \end{array}$$

Analysing Fig. 2 in the same manner as Fig. 1, we can say that in binary mixtures the association probability of the longer chain molecules,  $P_{22}$ , decreases as the chain length of the second component increases and  $N_1$  is constant. A similar table to  $P_{11}$  may be written for  $P_{22}$ , but with the reversed sens of inequalities:

$$\begin{array}{c}
 P_{22}(8,8/8,10) > P_{22}(8,8/8,12) > P_{22}(8,8/8,14) > \dots > P_{22}(8,8/8,24) \\
 \wedge \qquad \qquad \qquad \wedge \qquad \qquad \qquad \wedge \\
 P_{22}(10,10/10,12) > P_{22}(10,10/10,14) > \dots > P_{22}(10,10/10,24) \\
 \wedge \qquad \qquad \qquad \wedge \\
 P_{22}(12,12/12,14) > \dots > P_{22}(12,12/12,24) \\
 \wedge \\
 \dots \\
 \wedge \\
 (P_{22}(22,22/22,24))
 \end{array}$$

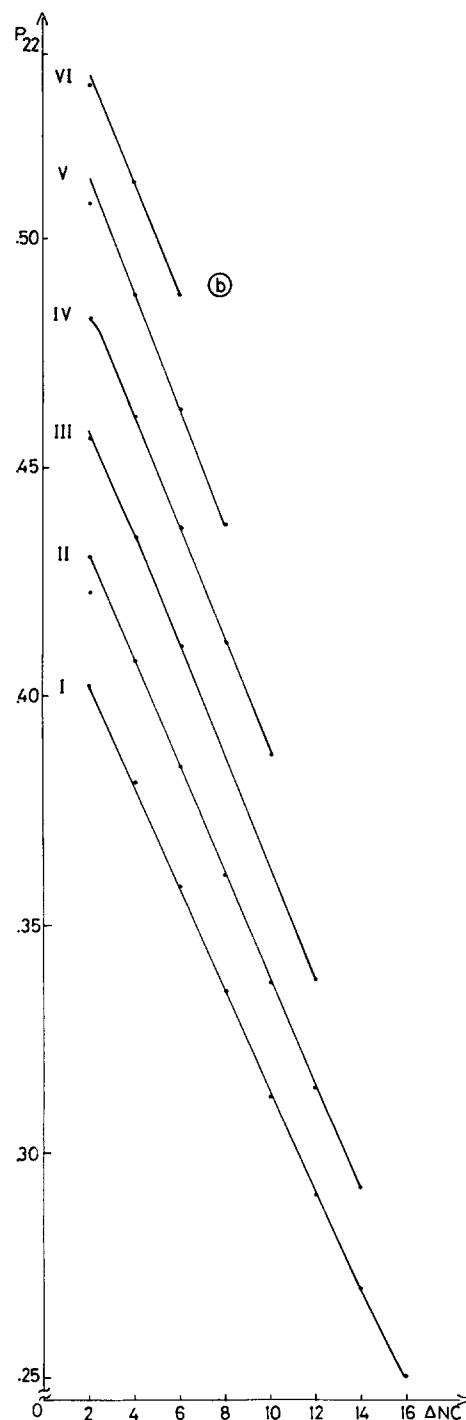
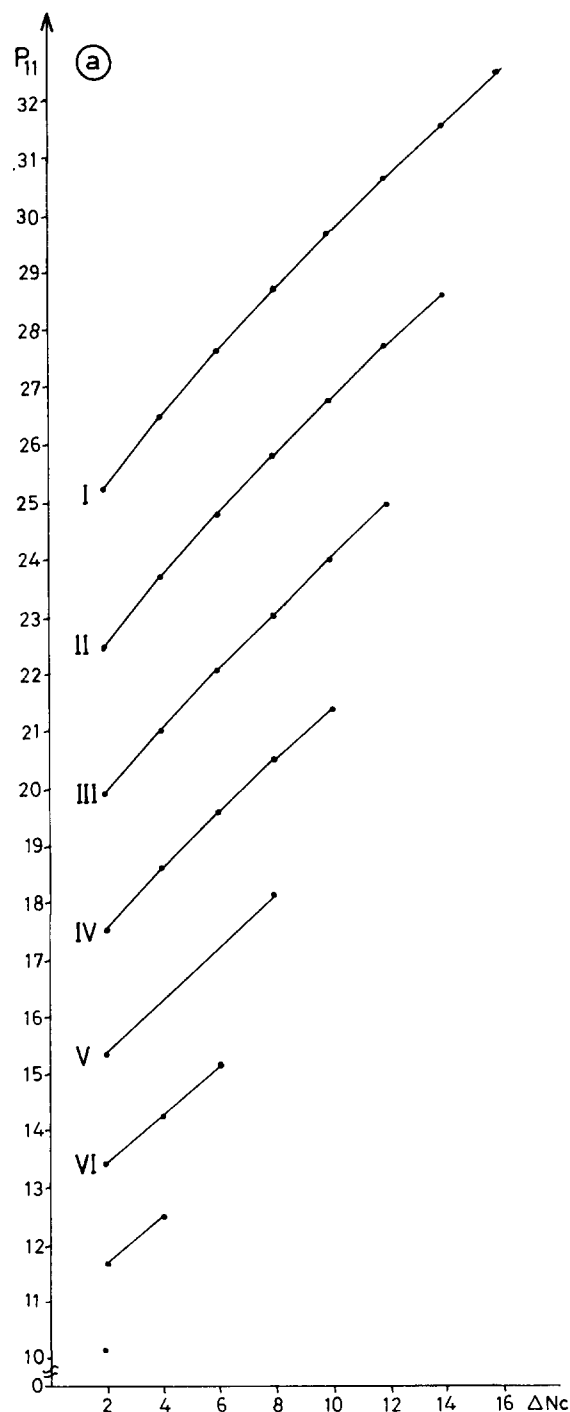


Fig. 4. The association probabilities variation  $P_{11}$  (a),  $P_{22}$  (b) and  $P_{12}$  (c) versus the difference between the lengths of the two mixed molecules species,  $\Delta N_c = N_2 - N_1$ , corresponding to each series. In (c) the curves for the other series which are situated (in increasing order according to the rank of the series) between the straight lines represented in the diagram. All curves have been drawn for  $x_1 = x_2 = 0.5$  and  $p_1 = p_2 = 5$  D.

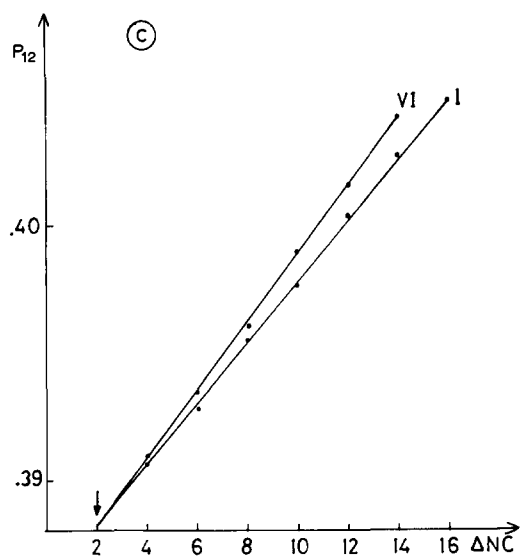


Fig. 4 (continued).

For both tables the sequences of inequalities are available for all values of the dipole moment of the second component of the mixture.

The cross association probability  $P_{12}$  between two amphiphiles with different chains increases as the chain length of the second increases ( $N_1 = \text{constant}$ ) and de-

creases when the chain length of the first component increases ( $N_2 = \text{constant}$ ) for each value of the dipole moment  $p_2$ .

All curves are nearly parallel in the Figs. 1 and 2. But, unlike the curves in Figs. 1 and 2, the slope of curves  $P_{12} = f(p_2)$  increases as the chain length of the head of series increases and the higher the dipole moment  $p_2$ , the higher the slope is. For this reason the curves corresponding to binary mixtures in which the second component is the same, get near one to another, when the dipole moment  $p_2$  increases. A symmetry is noticed to the mixtures in which the head of the series would have a chain with 15 carbon atoms. This is more clear for a high electric dipole moment ( $p_2 = 45$  D) when there the following approximate equalities:

$$P_{12}(8,24/8,24) \approx P_{12}(22,24/22,24) \quad (\text{Fig. 3a})$$

$$P_{12}(10,10/10,24) \approx P_{12}(20,20/20,24) \quad (\text{Fig. 3b})$$

$$P_{12}(10,22/10,22) \approx P_{12}(20,22/20,22)$$

$$P_{12}(12,24/12,24) \approx P_{12}(18,24/18,24)$$

$$P_{12}(12,22/12,22) \approx P_{12}(18,22/18,22) \quad (\text{Fig. 3c})$$

$$P_{12}(12,20/12,20) \approx P_{12}(18,20/18,20)$$

$$P_{12}(14,24/14,24) \approx P_{12}(16,24/16,24)$$

$$P_{12}(14,22/14,22) \approx P_{12}(16,22/16,22)$$

$$P_{12}(14,20/14,20) \approx P_{12}(16,20/16,20) \quad (\text{Fig. 3d})$$

$$P_{12}(14,18/14,18) \approx P_{12}(16,18/16,18)$$

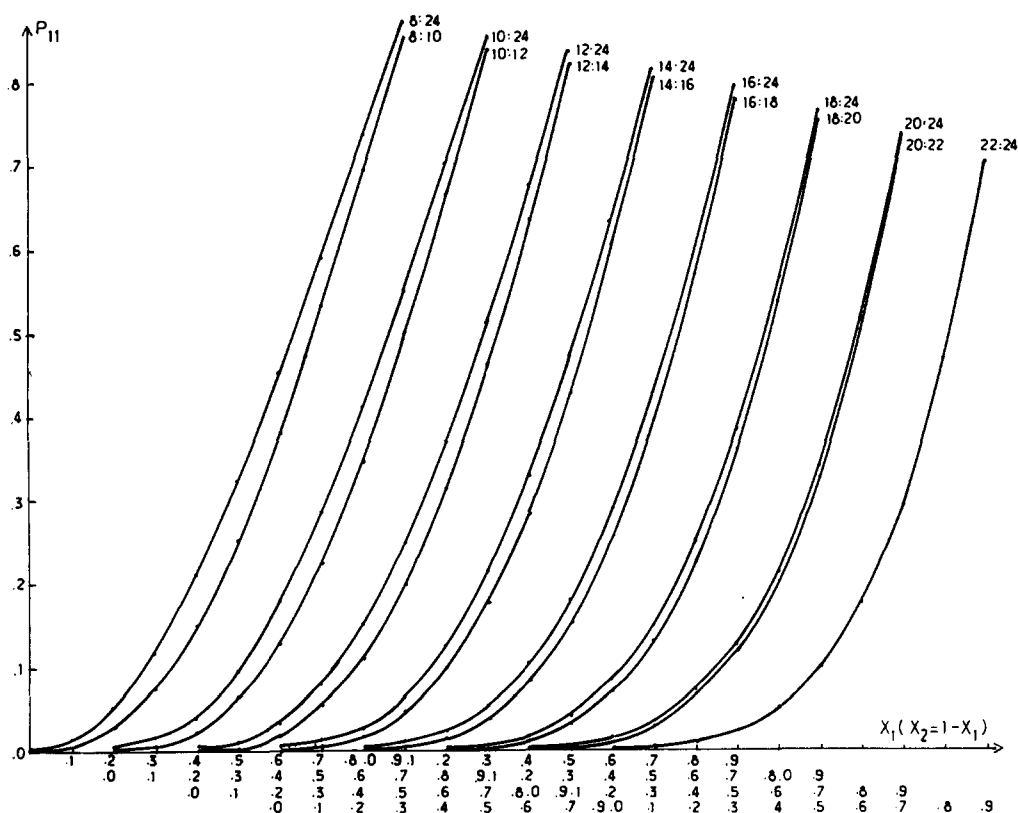


Fig. 5. The association probability variation  $P_{11}$  of the shorter chain molecules in a binary mixture organized in a planar bilayer depending on the molal fraction of the first component.  $p_1 = p_2 = 5$  D.



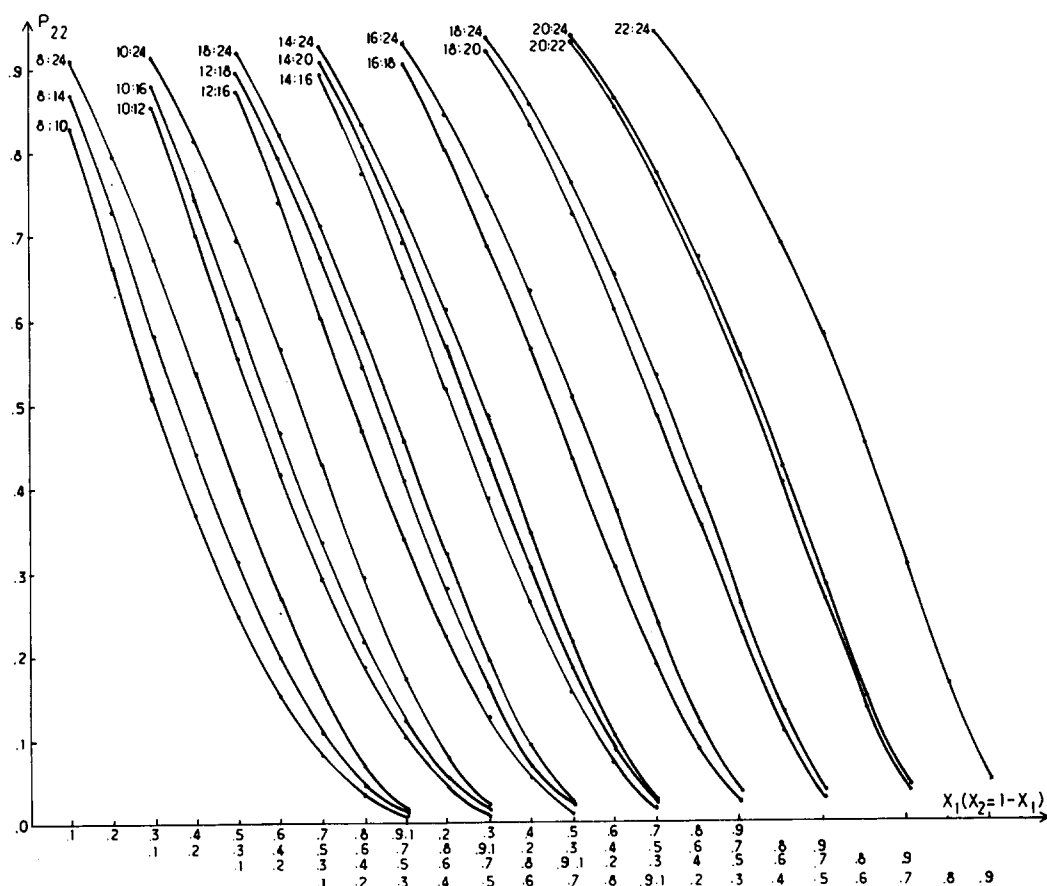


Fig. 6. The association probability variation  $P_{22}$  of the longer chain molecules in a binary mixture depending on the molal fraction of the first component. In Figs. 5 and 6 in every group of curves the order from bottom to top corresponds to the increase in the length of the second component. The order from the left to the right of the sets of curves corresponds to the increase in the hydrocarbon chain length of the head of series. The omitted curves are placed between the curves which was drawn.

One can write a table of inequalities for  $P_{12}$ , but the vertical sequences of inequalities will not be real for all values of the dipole moment  $p_2$ . We can sum up the conclusions resulted from Figs. 1–3 in a scheme valid for each value of the dipole moment  $p_2$  (Scheme I).

$N_2 = \text{constant}$  and  $N_1$  increases  $\Rightarrow P_{11}$  decreases;  $P_{12}$  decreases;  $P_{22}$  increases  
 $N_1 = \text{constant}$  and  $N_2$  increases  $\Rightarrow P_{11}$  increases;  $P_{12}$  increases;  $P_{22}$  decreases

Scheme I

This scheme shows a most interesting aspect: the insertion of a  $\text{CH}_2$  group into the shorter chain or into the longer chain has opposite effects on each of the three association probabilities,  $P_{11}$ ,  $P_{12}$  and  $P_{22}$ . At the same time if all the curves from Fig. 2 would be drawn on a single rectangular axes then some groupings of probability curves  $P_{22}$  from different series would appear. For example, the curves marked with the number pairs 8:12, 10:16, 12:20, 14:24 would form a group in this order.

For the binary systems corresponding to these curves grouped together, the association probability  $P_{22}$  is approximately the same for each value of the dipole moment  $p_2$ . For each group of such curves the short chain of the first component increases with two carbon atoms and the long chain of the second component increases with four carbon atoms. Taking into account that  $P_{22}$  increases when  $N_1$  increases and decreases when  $N_2$  increases, we can say that a carbon atom added to the short chain is twice more effective on the association probability  $P_{22}$  than a carbon atom added to the long chain.

Making the same reasoning for the association probability  $P_{11}$  we find that a carbon atom inserted into the short chain is nearly four times more effective than a carbon atom inserted into the long chain. For example, the following curves 12:16 and 14:24 will be nearly overlapped.

In each series of mixtures association probabilities  $P_{11}$  and  $P_{12}$  have a linear increasing variations (Figs. 4a,c) and  $P_{22}$  has a linear decreasing variation (Fig. 4b) as the difference between the lengths of the chains of

the two components of the mixture,  $\Delta N_c$  ( $\Delta N_c = N_1 - N_2$ ,  $N_1$  is constant for each series), increases.

The monotony of the straight lines from Figs. 4a,b,c confirms the conclusion from the second row of Scheme I, but at the same time, the shifting of curves in Figs. 4a,b confirms the difference between the efficiency of a carbon atom on  $P_{11}$  and  $P_{22}$ , when it is added to a shorter or a longer chain. So, adding a carbon atom to the short chain the diminishing of  $P_{11}$  is greater than the increasing of  $P_{11}$  due to the adding of a carbon atom to the long chain, resulting, a shift of lines towards the bottom of Fig. 4a. Also, adding a carbon atom to the short chain, the increase of  $P_{22}$  is higher than the decrease of  $P_{22}$  due to the adding of a carbon

atom to the long chain, resulting a shift of lines towards the top of Fig. 4b.

If we introduce the notion of the ratio of efficiency,  $k$ , as the ratio between the number of carbon atoms added to the short chain and the number of carbon atoms added to the longer chain that modifies the association probability with the same quantity (absolute value) we can sum up that:

$$k = \begin{cases} 1/2 & \text{for the association probability } P_{11} \\ 1/4 & \text{for the association probability } P_{12} \\ 1/4 & \text{for the association probability } P_{22} \end{cases} \quad (11)$$

The latter value results from the fact that the lines which represent the  $P_{12}$  versus  $\Delta N_c$  do not move, but

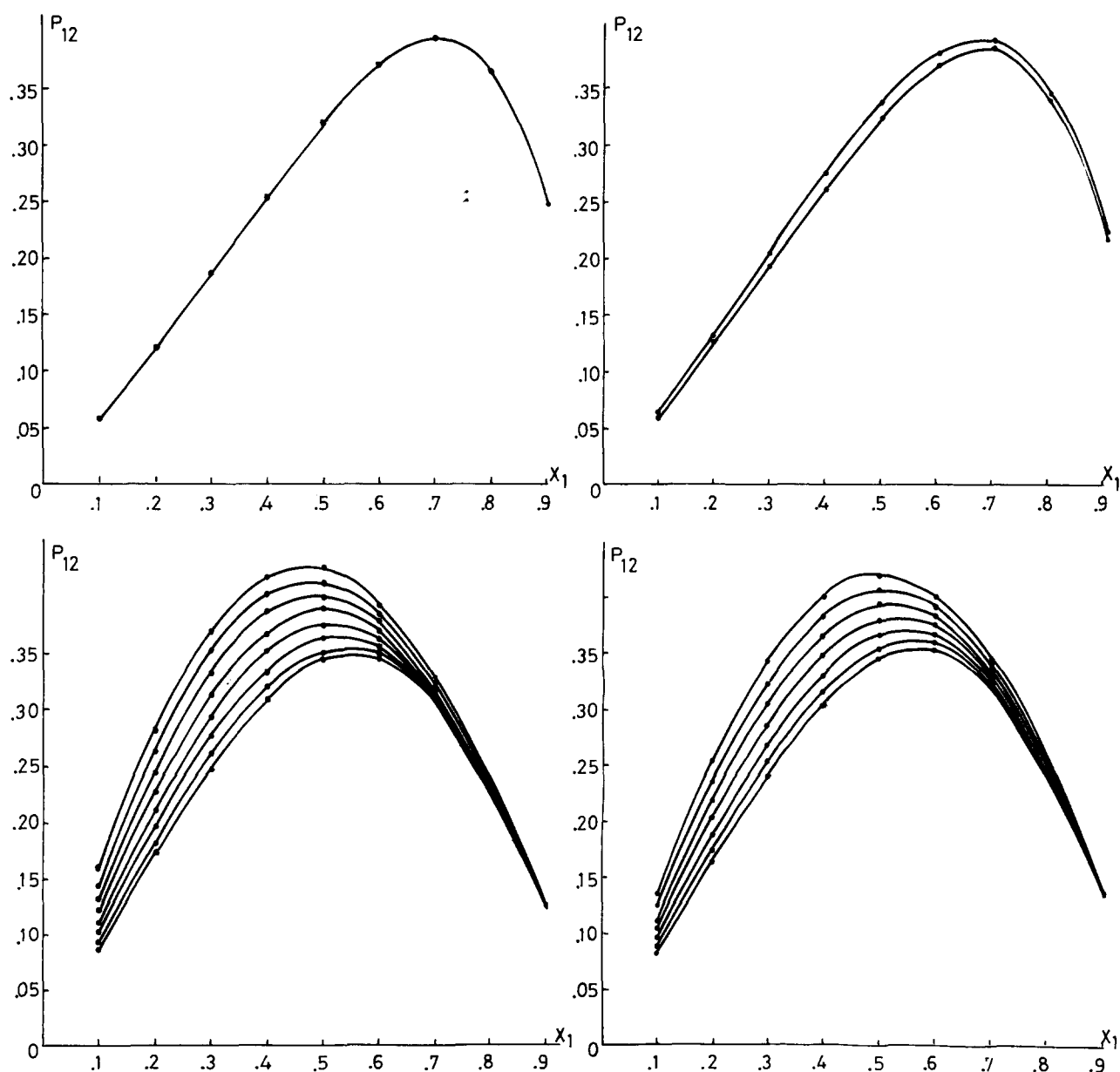


Fig. 7. The cross association probability variation  $P_{12}$  depending on the molal fraction of the first component. The series VIII, VII, I, II.

make a slight rotation (Fig. 4c). From this rotation and owing to the straight lines in Figs. 4a,b there are approximately parallel results that the ratio of efficiency depends slowly on the difference of chain lengths.

Normally, if the shorter chain molecules concentration increases, then the association probability  $P_{11}$  increases (Fig. 5), the association probability  $P_{22}$  decreases (Fig. 6) and  $P_{12}$  has a parabolic variation whose maximum increases and moves towards the small concentrations of the first component when the chain length of the second component increases (Figs. 7 and 8).

It is useful not only to know how the association probabilities  $P_{11}$ ,  $P_{12}$ ,  $P_{22}$  depend on each of the

parameters ( $N_1$ ,  $N_2$ ,  $p_2$ ,  $x_1$ ) but also their absolute values and their succession for some combination of values of these parameters. Fig. 9 gives all the binary mixtures that have been studied. The numerator is the shorter chain and the denominator is the longer chain of the molecules in the binary mixture. The solid broken line divides all the binary mixtures in three classes, each of them being characterised by a certain succession of the three association probabilities, available when  $p_1 = p_2 = 5 D$ , and  $x_1 = x_2 = 0.5$ . Each class contains all binary mixtures for which one of the following succession of inequalities is valid:  $P_{22} < P_{11} < P_{12}$  (class A);  $P_{11} < P_{22} < P_{12}$  (class B) and  $P_{11} < P_{12} < P_{22}$  (class C).

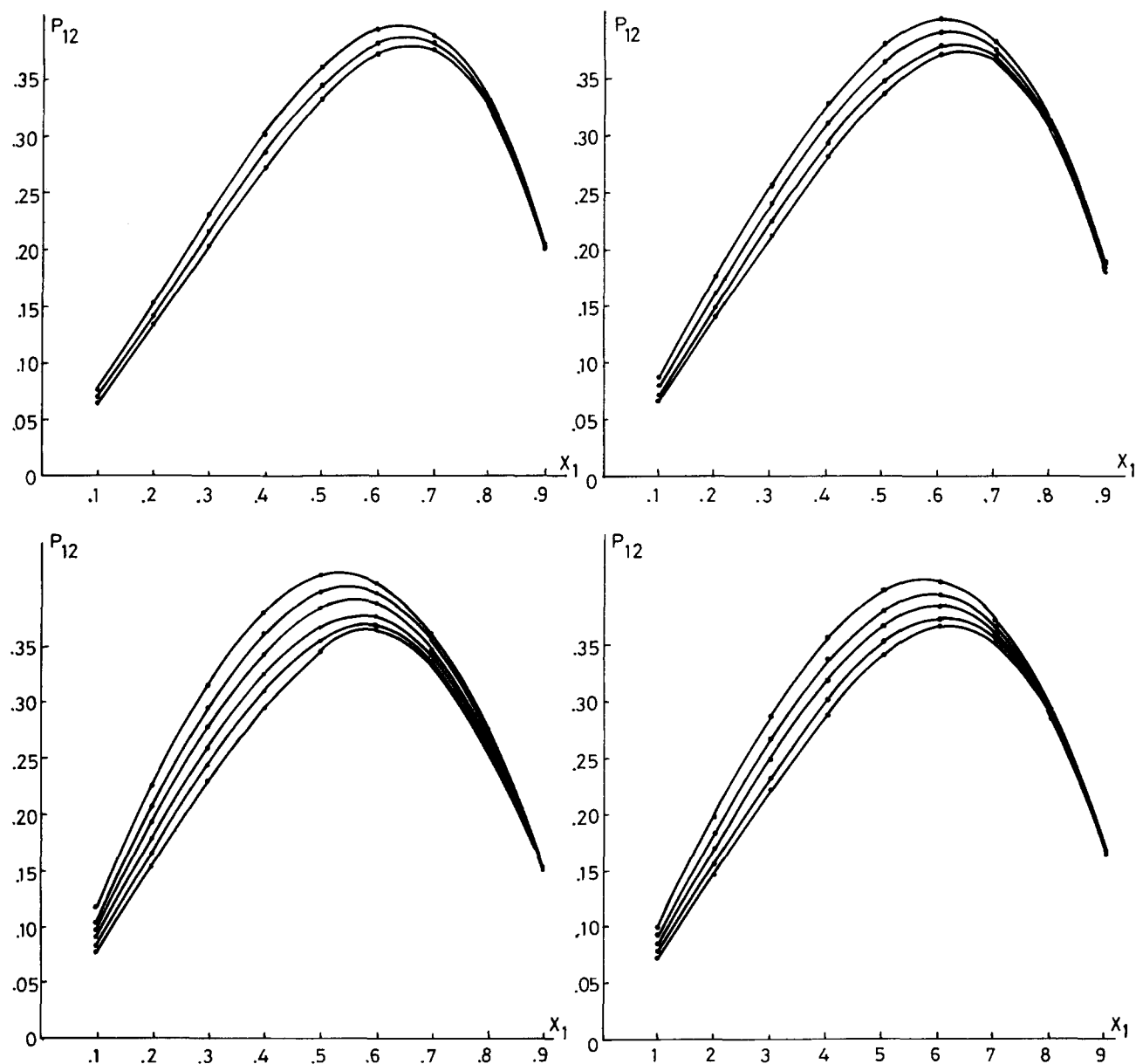


Fig. 8. The cross association probability variation  $P_{12}$  depending on the molal fraction of the first component. The series VI, V, III, IV. In Figs. 7 and 8 the reading of the series one make from left to the right, beginning at the top of the picture. In each group of curves the order from bottom to top corresponds to the increase in the length of the second component.  $x_1 = x_2 = 0.5$  and  $p_1 = p_2 = 5 D$ .

With the increase of the electric dipole moment  $p_2$  some binary mixtures pass from a class into another. In Fig. 9 the circled number to the left of each fraction represents the value of the dipole moment for which this binary mixture passes from class A into class B, and the circled number to the right of each fraction is the value of the dipole moment for which this binary mixture passes from class B into class C. For example the dashed broken line points the division of all binary mixtures into the three classes when  $p_1 = 5$  D and  $p_2 = 35$  D. Finally (when  $p_1 = 5$  D and  $p_2 = 45$  D) there are only two classes, B and C, separated by the line marked with  $A_2$  in Fig. 9. So, the cross association of single chain molecules is favoured for higher electric dipole moment.

It is well known that the hydrophobic chains of natural phospholipids are unequal, having 14–22 carbon atoms [5] and the dipole moment  $p \geq 35$  D [6]. It is a very interesting coincidence the fact that for the binary systems formed of single chain amphiphile molecules that have the chain length of 14 to 22 carbon atoms,  $p_1 = 5$  D,  $p_2 \geq 35$  D and  $x_1 = x_2 = 0.5$ , the cross association probability  $P_{12}$  is much higher than the other two ( $P_{12} > P_{22} > P_{11}$ ). On the other hand, even both chains of a phospholipid molecule have equal lengths the molecules adopt such a conformation that their chains appear to be unequal. This has been observed in crystal-line state of the phospholipid bilayers using X-rays

diffraction [7]. Therefore, the selective association of single chain molecules in a binary mixture could explain some problems regarding the aggregation of the lipid molecules into supermolecular structures: either micelles or plane and spherical bilayers.

We think that the stability of such supermolecular structures increases if  $P_{12}$  increases, because these aggregates would have an ordered structure of 'cells' (like crystal) each 'cell' being formed of a cross pair of unequal chain molecules. Accordingly, a bilayer can be more stable because of the possible interdigitation of the hydrophobic unequal chains from the two monolayers of a bilayer. The micelles formed of mixture of single unequal molecules should be more stable, too.

A lot of binary mixtures studied here have the cross association probability  $P_{12}$  larger than  $P_{11}$  or  $P_{22}$ . As the electric dipole moment  $p_2$  increases the number of binary systems for which  $P_{12}$  is larger than  $P_{11}$  and  $P_{22}$  increases and  $P_{12}$  increases too.

On the other hand, the selective association between molecules can be the cause of phase separations in binary or ternary mixtures making up separate clusters through microinterfaces [1]. These preferential association give birth to structural nonhomogeneities with different degree of stability and are the first cause of thickness fluctuations of the bilayers. The separation of lipids has been detected in binary lipid mixtures in the gel phase using differential calorimetry with scanning

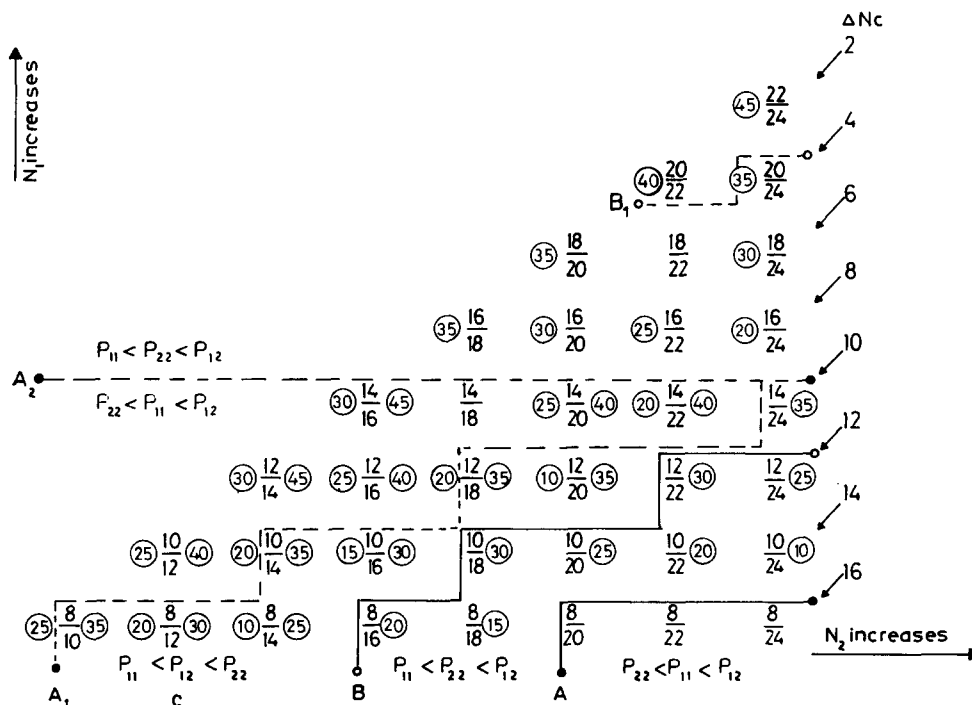


Fig. 9. The diagram of the all binary mixtures studied in this work. Under the diagram there are the inequalities between the association probabilities valid for each class. This diagram is useful to know the composition of each class for different values of the electric dipole moment  $p_2$ .

[8]. There are experimental attempts of studying liquid crystals made up of such biological lipid mixtures [9,10].

When one of the association probabilities is much smaller than the other two, these clusters are like little defects in the aggregate structure.

The values of  $a_1$  and  $a_2$  influence both dipole-dipole interactions and disperse forces, but detailed effects on the association probabilities will be given in a future paper.

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